



Combined Quantum Mechanics and Classical Electrodynamics Multiscale Approach for the Calculation of SERS Spectra (A Brief Survey)

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- **Key properties of *metallic* nanostructures: possibility of *collective excitation* of the *conduction* electrons by UV/Vis light**
- **These *surface plasmon excitations* are responsible for remarkable size / shape / environment-dependent *optical properties***
- **Characterization of metallic nanoparticles in combination with detailed quantitative *electromagnetic (EM) simulations* enabled synthesis of particles with pre-determined *spectral properties***
- **This control of the *optical* properties of nanomaterials resulted in a *wide range of applications* in ultra-sensitive chemical and biological sensing**



- Plasmon excitations lead to ***strongly enhanced EM fields*** near the nanoparticle's surface
- This f.i. results in intense ***absorption, fluorescence*** and ***scattering*** characteristics of the nanoparticles,
- And is responsible for the ***EM contribution*** to the enhanced Raman signals (up to $\sim 10^{10}$, single molecule spectroscopy) observed in ***surface-enhanced Raman scattering (SERS)***
- A ***complete*** picture of the various ***enhancement mechanisms*** (see below) is not available, due to the highly ***complicated experimental conditions*** (f.i. roughened surfaces, nanoparticle aggregates, chemical interactions of adsorbants to surfaces)



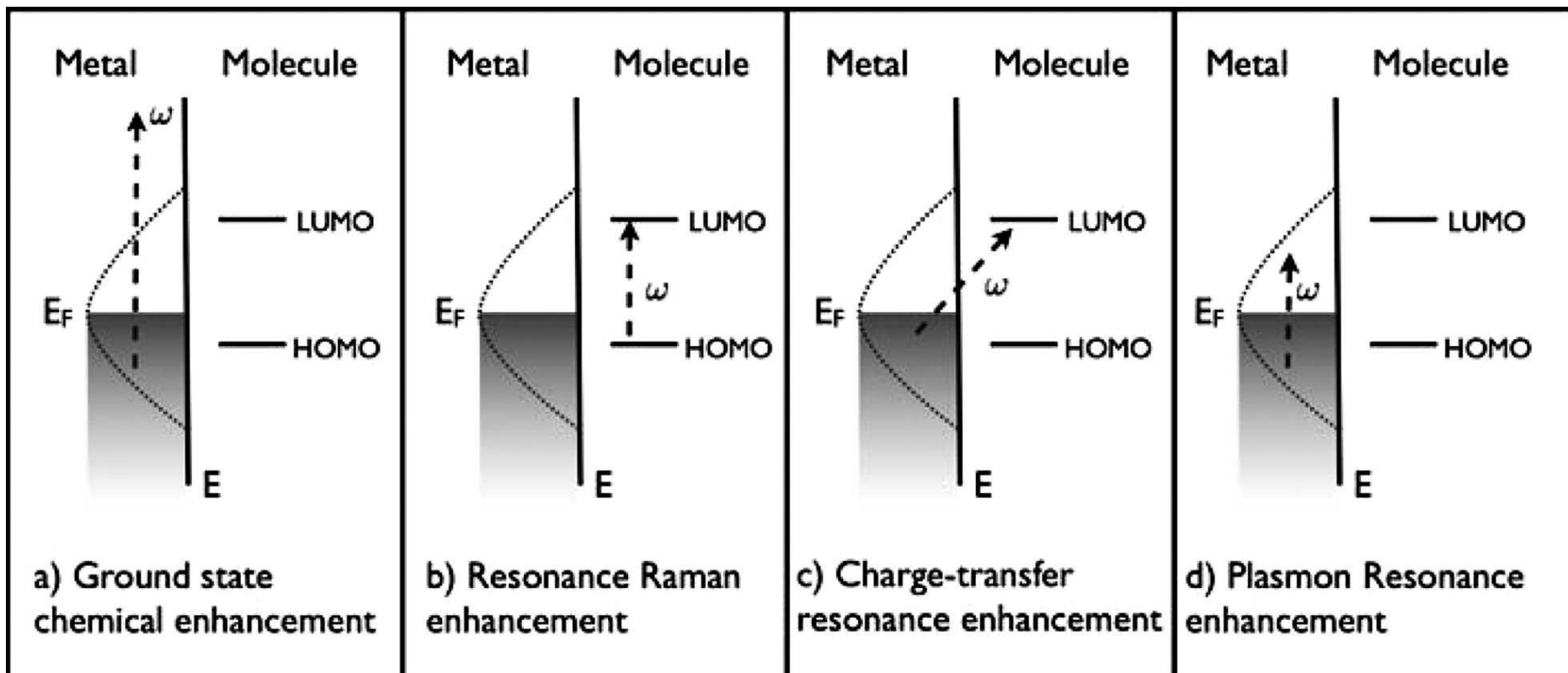
- Much progress has been achieved by accurate and efficient classical *electrodynamic (ED) simulations* of nanostructure optical properties
- Numerical methods such as f.i. *finite-difference time-domain (FDTD, see below)* approaches were used to calculate the *plasmonic properties* of complex shapes and arrangements
- These *computational* procedures lead to *detailed* insight into the **EM** mechanism in SERS, but they do *not* provide any informations with respect to *chemical enhancement (CHEM)* (*molecules* are treated as *dipoles* or are *neglected* at all!)



- **Challenge** in building a theory for calculating **SERS optical response**, that includes both **quantum mechanics (QM)** and **classical electrodynamics (ED)**:
 - Bridging the **length scales** needed for both approaches, which differ in **order of magnitude!**
- **Purely chemical** models of SERS based on **QM** are generally limited to **~1 nm** in size, **including** the metal particle/cluster
- **EM field evaluations** are usually based on **grids** or **finite elements**, that have **1 nm** dimensions at the **minimum**



- ***For this reason***, theoretical treatments of SERS often take ***one of two*** paths:
 - One approach ***neglects*** the **CHEM** enhancement and ***focuses*** on the ***predominant EM*** enhancement
 - Other studies ***only determine*** the **CHEM** enhancement using ***small atomic cluster*** models of the nanoparticle
- Much of ***current*** research in this field focuses on ***novel multiscale approaches*** for analysis and understanding SERS ***mechanisms*** by
 - Combination of ***quantum mechanics*** (e.g. **RT-TDDFT**, **LR-TDDFT**) and ***classical electrodynamics*** (f.i. **FDTD**, **FEM**, **MMP**) methods



(QM)

(QM)

(QM)

(ED+QM)

(According to G.C. Schatz et al., 2008)



- a) Enhancement due to ground state *chemical interactions (CHEM)* between molecule (adsorbate) and nanoparticle/surface, not associated with any *electronic excitations* of the nanoparticle-molecule system, λ_{exc} arbitrarily chosen (UV/Vis, IR laser) (*non-resonant*)

- b) *Resonance Raman (RR)* enhancement with λ_{exc} being *resonant* with a *molecular electronic* transition

- c) *Charge-transfer (CT)* resonance Raman enhancement with λ_{exc} being *resonant* to nanoparticle-adsorbant *CT-transitions*

- d) Enhancement due to a *very strong local field (EM)*, when λ_{exc} is *resonant* with *the plasmon excitations* in the metal nanoparticle



- ***Raman scattering intensity*** for the free molecule is given by:

$$I_M^R \propto \left| \frac{\partial \alpha_M}{\partial Q_M} \right|^2$$

α_M : molecular polarizability

Q_M : a normal mode of the molecule



- Raman scattering of the molecule *is affected* by **EM** interaction with a *polarizable* body (metallic particle) under radiation, *located close* to the molecule
- The Raman intensity can *now* be *expressed* as:

$$I^R = I_M^R \cdot |E^{loc}|^4$$

E^{loc} : *local field enhancement* due to metal nanoparticle, get that from *classical ED simulations*, f.i. **FDTD** (!)



- ❖ Light is assumed to *incident* on a system, that is *discretized* into *many small* buildings blocks
- ❖ Each of them is characterized by a *dielectric permittivity*, $\epsilon(\vec{r})$, and by a *magnetic permeability*, $\mu(\vec{r})$ (material's properties)
- ❖ Then *Maxwell's equations* (see below) are solved in the *real time domain* to obtain the *time evolution* of
 - the *electric* field, $\vec{E}(\vec{r}, t)$,
 - the *magnetic* field, $\vec{B}(\vec{r}, t)$, and
 - the *electric current density* $\vec{J}(\vec{r}, t)$



Name	<u>Integral</u> equations	<u>Differential</u> equations
<u>Gauss's law</u>	$\oiint_{\partial\Omega} \mathbf{E} \cdot d\mathbf{S} = \frac{1}{\epsilon_0} \iiint_{\Omega} \rho dV$	$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$
<u>Gauss's law for magnetism</u>	$\oiint_{\partial\Omega} \mathbf{B} \cdot d\mathbf{S} = 0$	$\nabla \cdot \mathbf{B} = 0$
Maxwell–Faraday equation (<u>Faraday's law of induction</u>)	$\oint_{\partial\Sigma} \mathbf{E} \cdot d\boldsymbol{\ell} = -\frac{d}{dt} \iint_{\Sigma} \mathbf{B} \cdot d\mathbf{S}$	$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$
<u>Ampère's circuital law</u> (with Maxwell's correction)	$\oint_{\partial\Sigma} \mathbf{B} \cdot d\boldsymbol{\ell} = \mu_0 \iint_{\Sigma} \left(\mathbf{J} + \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \right) \cdot d\mathbf{S}$	$\nabla \times \mathbf{B} = \mu_0 \left(\mathbf{J} + \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \right)$



$$\begin{aligned}\nabla \cdot \mathbf{E} &= 0 & \nabla \times \mathbf{E} &= -\frac{\partial \mathbf{B}}{\partial t}, \\ \nabla \cdot \mathbf{B} &= 0 & \nabla \times \mathbf{B} &= \frac{1}{c^2} \frac{\partial \mathbf{E}}{\partial t}.\end{aligned}$$

$$\frac{1}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2} - \nabla^2 \mathbf{E} = 0, \quad \frac{1}{c^2} \frac{\partial^2 \mathbf{B}}{\partial t^2} - \nabla^2 \mathbf{B} = 0,$$

$$c = \frac{1}{\sqrt{\mu_0 \epsilon_0}} = 2.99792458 \times 10^8 \text{ m s}^{-1}$$

$$c = \frac{1}{\sqrt{\mu_0 \mu_r \epsilon_0 \epsilon_r}} \quad (\text{Materials})$$



- ❖ The **EM properties** can also be determined in the **frequency domain** through **Fourier transform**
- ❖ The **total electric field**, $\vec{E}_{total}(\vec{r}, \omega)$, at a given observation point, is then a **combination** of the **scattered field**, $\vec{E}_{sca}(\vec{r}, \omega)$, and of the **incident field**, $\vec{E}_0(\vec{r}, \omega)$:

$$\vec{E}_{total}(\vec{r}, \omega) = \vec{E}_{sca}(\vec{r}, \omega) + \vec{E}_0(\vec{r}, \omega)$$

- ❖ A **scattering response function**, $\lambda(\vec{r}, \omega)$, which is defined as

$$\lambda(\vec{r}, \omega) \sim \frac{\vec{E}_{sca}(\vec{r}, \omega)}{\vec{E}_0(\vec{r}, \omega)},$$

provides a measure of the **local field enhancement**



- **Dynamic (frequency-dependent) polarizabilities are *necessary* for e.g.**
 - Calculation of **Resonance Raman Spectra (RRS)** or
 - **Hybrid Quantum Mechanics / Classical Electrodynamics simulations (QM/ED)**
- May ***be quantumchemically*** obtained by f.i.
 - **Real-time time-dependent density functional theory (RT-TDDFT)** (e.g. Octopus, NWChem)
 - **"Excited state gradient" (EG) or "Short-time approximation" (STA)** (f.i. Gaussian, TURBOMOLE, etc.)
 - **Time-dependent density functional theory (TDDFT, linear response) or "Polarizability method" (PM)** (e.g. ADF, Gaussian, NWChem)



- For a molecule exposed to a *time-dependent external electric* Field, E_i , along axis i , the *dipole moment*, P_j , along axis j , in linear *first-order* approximation, is

$$P_j = P_{j0} + \alpha_{ij} \cdot E_i ,$$

where P_{j0} is the *permanent* dipole moment and α_{ij} represents the linear *polarizability tensor*

- In the *time* domain, one may then write

$$P_j(t) = P_{j0} + \int \frac{d\omega}{2\pi} e^{-i\omega t} \alpha_{ij}(\omega) \cdot E_i(\omega) ,$$

and the *induced* dipole moment, $P_j^{Ind}(t)$, is *defined* as

$$P_j^{Ind}(t) = P_j(t) - P_{j0}$$



- In the *frequency* domain, we obtain

$$P_j^{Ind}(\omega) = \alpha_{ij}(\omega) \cdot E_i(\omega) ,$$

where

$$\alpha_{ij}(\omega) = \frac{P_j^{Ind}(\omega)}{E_i(\omega)} = \frac{\int dt e^{+i\omega t} P_j^{Ind}(t) e^{-\Gamma t}}{\int dt e^{+i\omega t} E_i(t)}$$

- This equation *relates* the *frequency-dependent polarizability tensor*, $\alpha_{ij}(\omega)$, to the *time evolution* of the molecule's *induced* dipole moment, $P_j^{Ind}(t)$, under a *time-dependent* external electric field, $E_i(t)$
- This procedure allows *incorporation* of the effect of *coupling* to the *metal particle* on the *excited state dynamics* of the molecule



- Within the framework of **DFT**, the *time-dependent* dipole moment, $\vec{P}(t)$, can be calculated from the *perturbed* electron density, which arises, when the system is *subjected* to an *applied electric* field, $\vec{E}_0(t)$
- The following *time-dependent* Schrödinger equation (TDSE), for this reason, is used:

$$i \frac{\partial}{\partial t} \varphi(r, t) = \left[-\frac{1}{2} \nabla^2 + \int dr' \frac{\rho(r', t)}{|r - r'|} + \frac{\delta E_{xc}[\rho(r, t)]}{\delta \rho(r, t)} - \vec{E}_0(t) \cdot \vec{r} \right] \varphi(r, t)$$

- The *coupling* Hamiltonian between the *external* electric field and the *molecule* is given by:

$$- \int \varphi^*(r) \vec{E}_0(t) \cdot \vec{r} \varphi(r) dr = -\vec{E}_0(t) \cdot \int \varphi^*(r) \vec{r} \varphi(r) dr = -\vec{E}_0(t) \cdot \vec{P}(r)$$

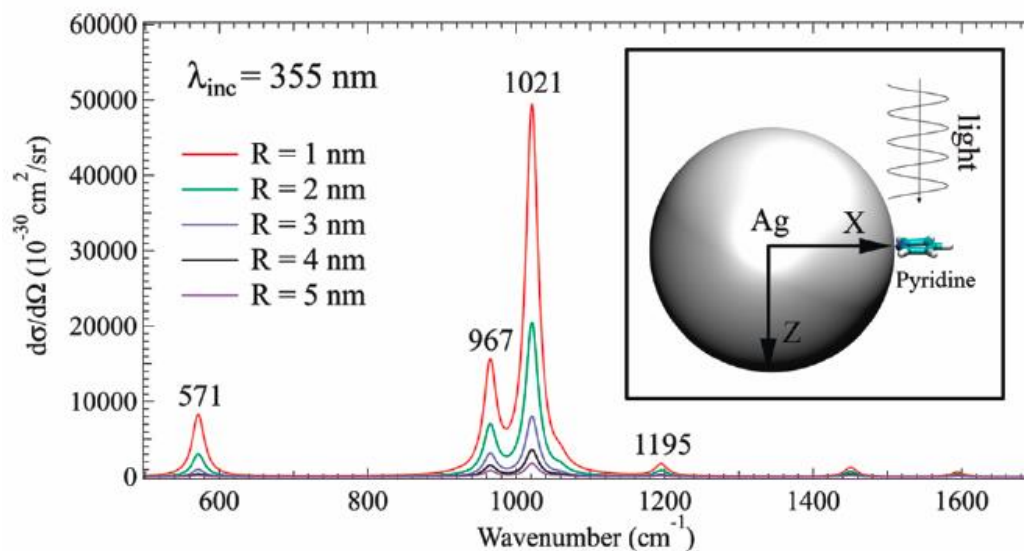
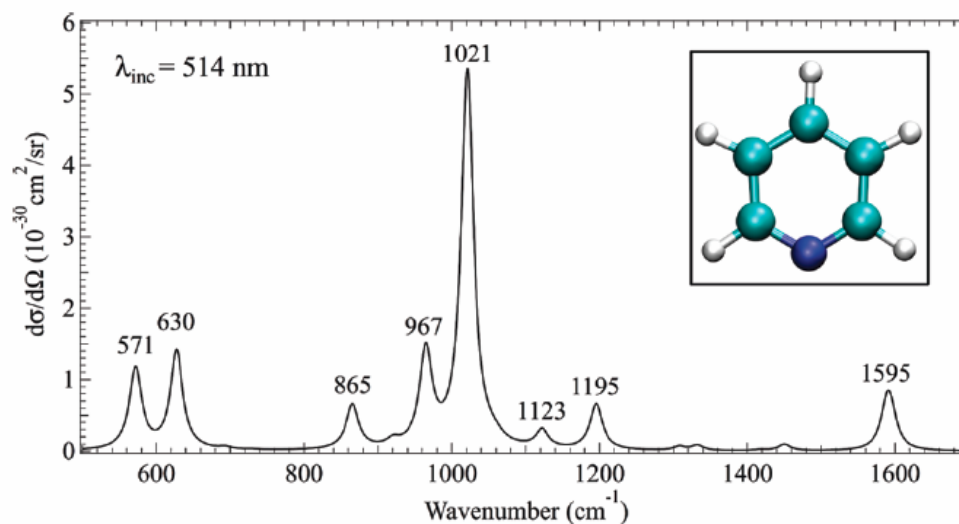


- Under the assumption of a *uniform scattered* electric field, $\vec{E}_{sca}(\vec{r}, t)$, the *Hamiltonian* of *the adsorbate molecule* in the *presence* of an incident field, $\vec{E}_0(\vec{r}, t)$, at the **RT-TDDFT** level of theory, can be written as:

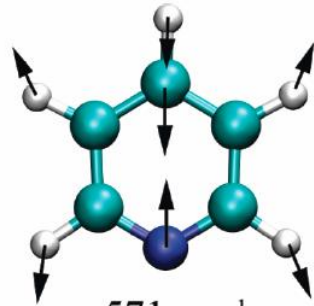
$$\hat{H}(t) = -\frac{1}{2} \nabla^2 + \int dr' \frac{\rho(r', t)}{|\mathbf{r} - \mathbf{r}'|} + \frac{\delta E_{xc}[\rho(\mathbf{r}, t)]}{\delta \rho(\mathbf{r}, t)} - \vec{E}_0(t) \cdot \vec{r} - \vec{E}_{sca}(t) \cdot \vec{r}$$

- The scattered field, *imposed* by the *polarized* nanoparticle, can be obtained via the above *scattering response function* $\lambda(\vec{r}, \omega)$

Standard example from the literature: Adsorption of pyridine on Silver (1)

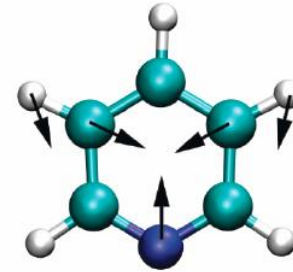


Standard example from the literature: Adsorption of pyridine on Silver (2)



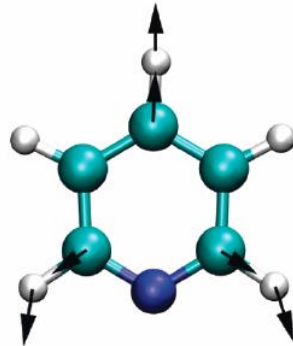
571 cm^{-1}

(a) ring deformation



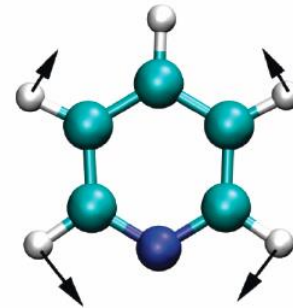
967 cm^{-1}

(b) ring breathing



1021 cm^{-1}

(c) ring breathing

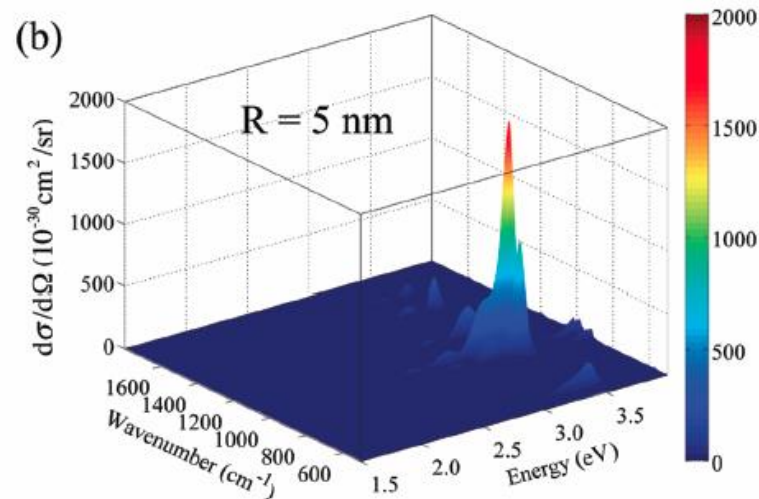
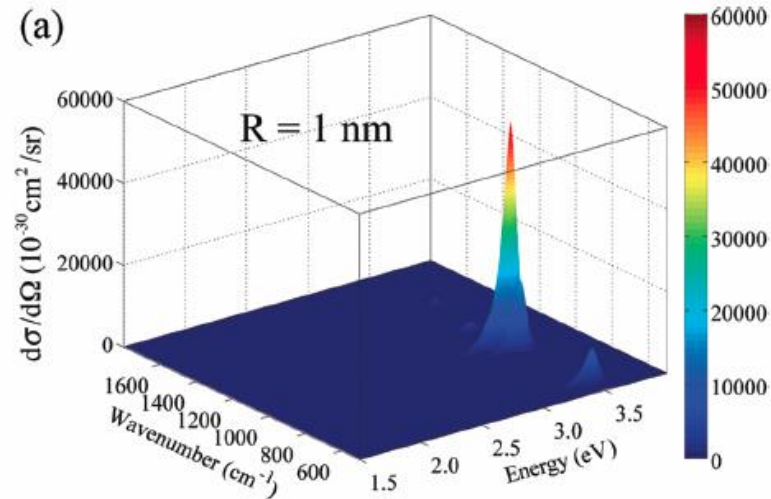


1195 cm^{-1}

(d) ring stretch

***Plasmon-enhanced* vibrational modes**

Standard example from the literature: Adsorption of pyridine on Silver (3)



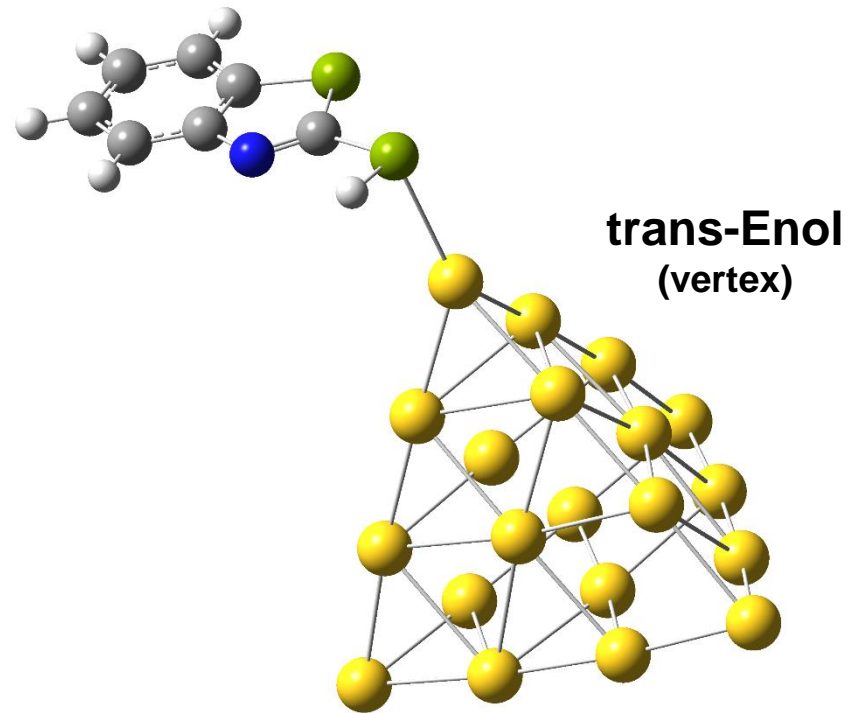
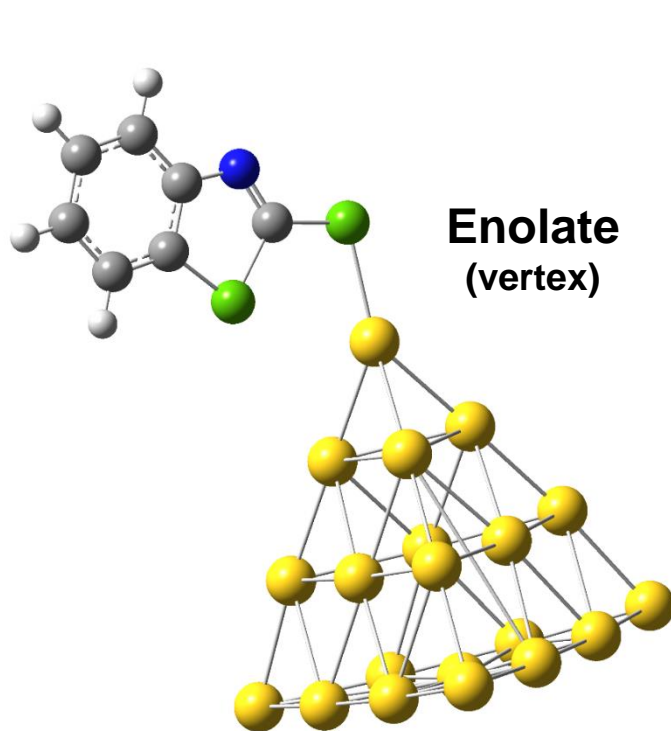
Surface plots of *differential Raman scattering cross section* as a function of incident light and vibrational mode wavenumber

Adsorption of MBT on gold (1)

(MBT = 2-Mercaptobenzothiazole)



- Computational model: **MBT(enolate/enol)** adsorbed on **tetrahedral Au₂₀**



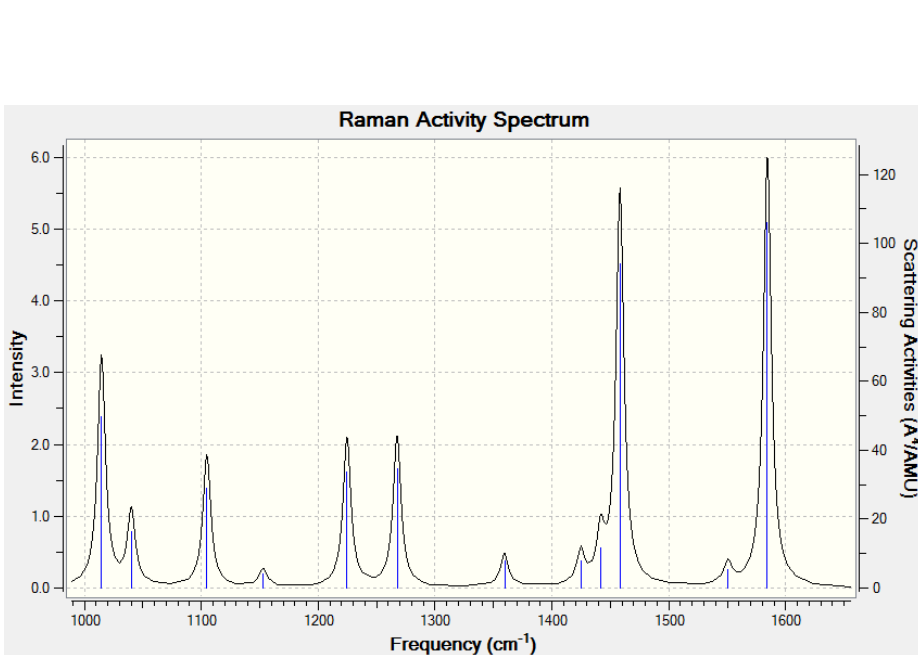
Different adsorption geometries!

Calculated relative stabilities for MBT(Enolate/Enol)-Au₂₀ (2)

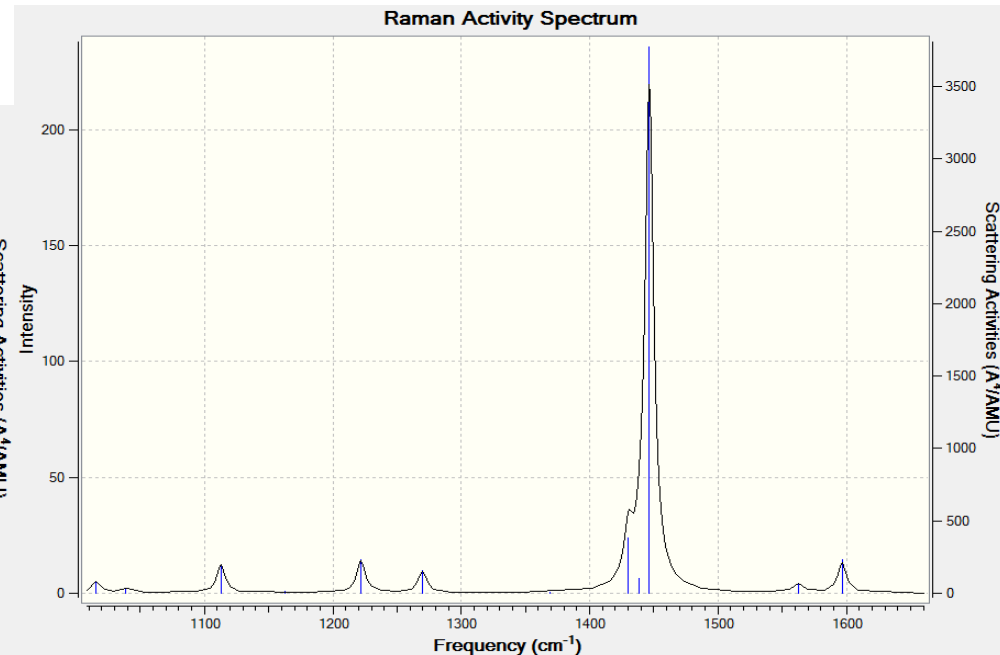


- Gaussian09:** Optimized geometries, Raman and IR vibrational frequencies and intensities, all structures considered represent minima on the respective BO energy hypersurface, no BSSE corrections applied for interaction energies

	Stabilization energies [kcal/mol]	
	Enolate (vertex) ("ad-atom" adsorption)	Enol (vertex) ("ad-atom" adsorption)
PBE/LanL2DZ (QZVP)	-53.2 (-56.8)	-9.2
B3LYP/LanL2DZ (QZVP)	-47.4 (-50.6)	-5.9
HCTH/LanL2DZ (QZVP)	-43.5 (-48.8)	-4.1
TPSS/LanL2DZ (QZVP)	-50.8 (-54.6)	-7.9
	"Chemisorption"	"Physisorption"

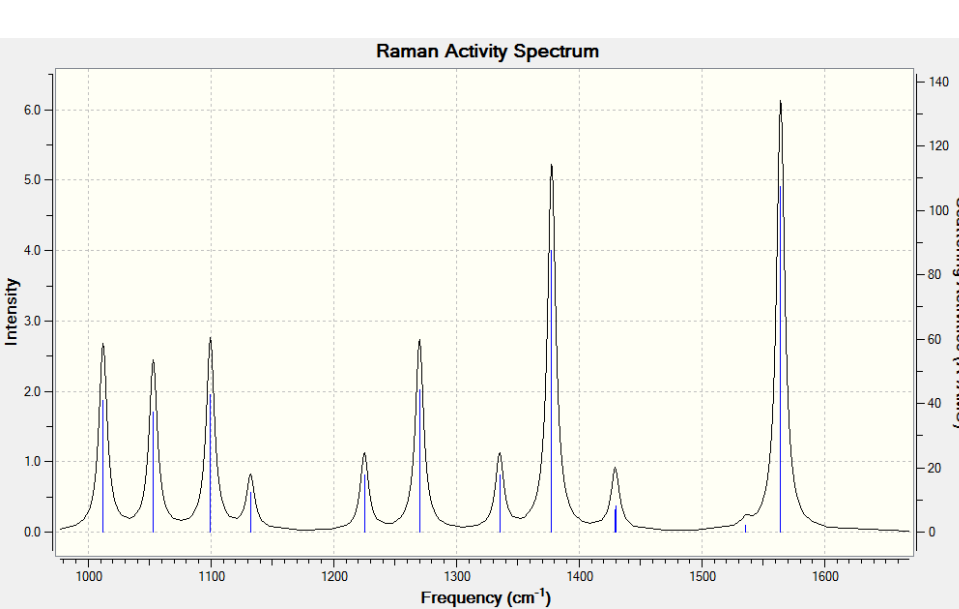


Free MBT molecule (Enolate)

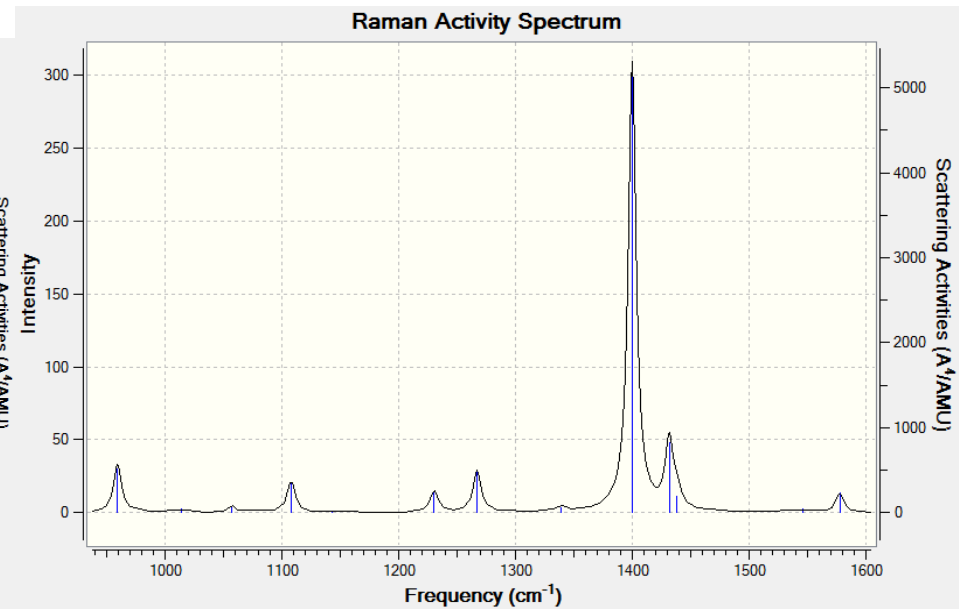


MBT(Enolate) on Au_{20}

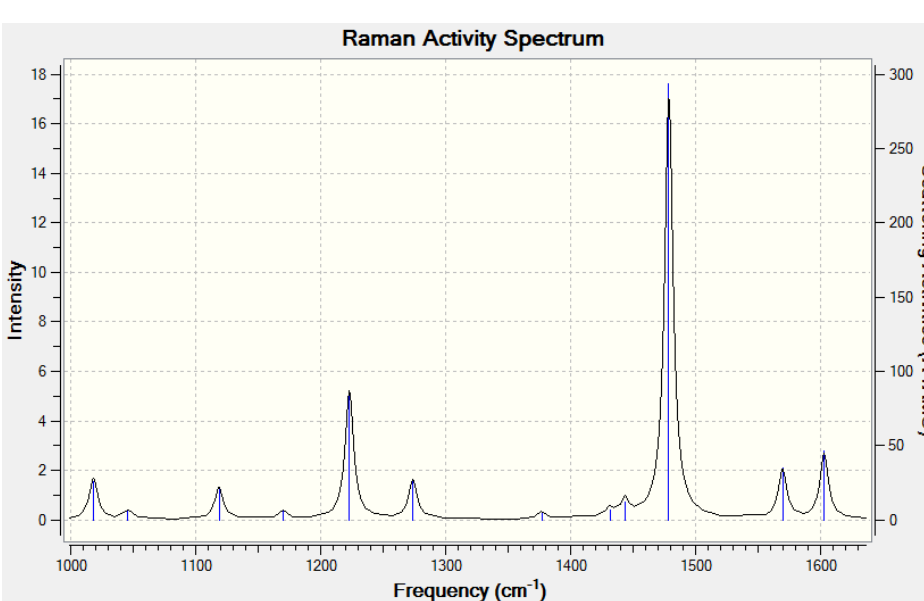
Calculated Raman spectra: "PBE/QZVP" (4)



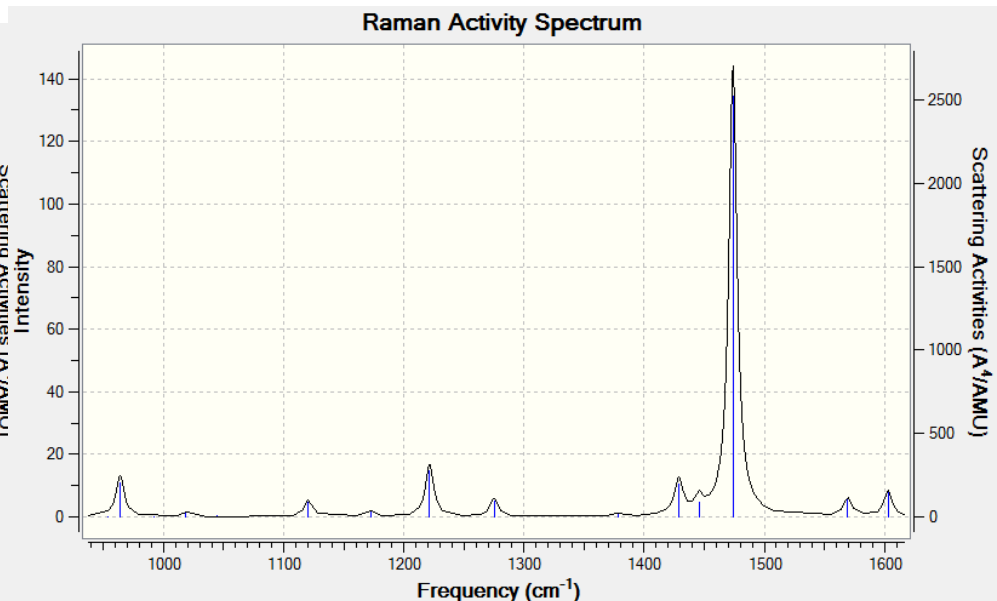
Free MBT molecule (Enolate)



MBT(Enolate) on Au_{20}



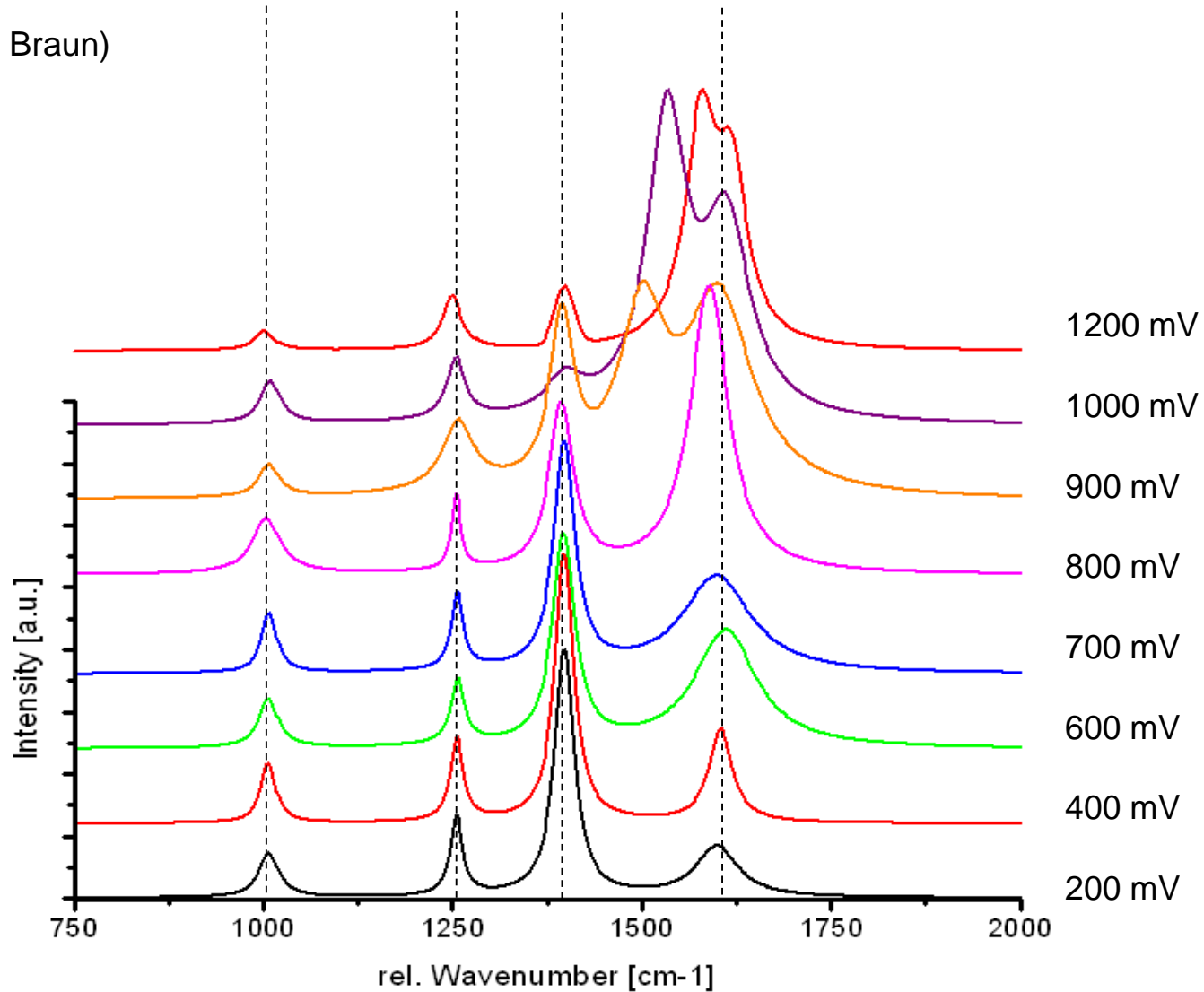
Free MBT molecule (trans-Enol)

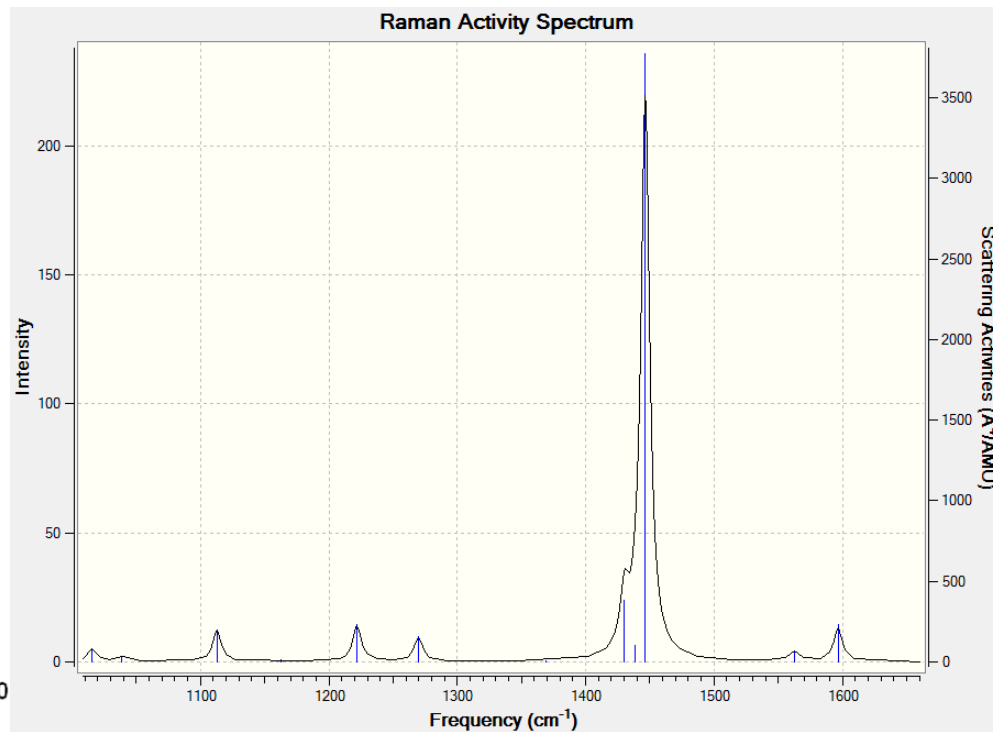
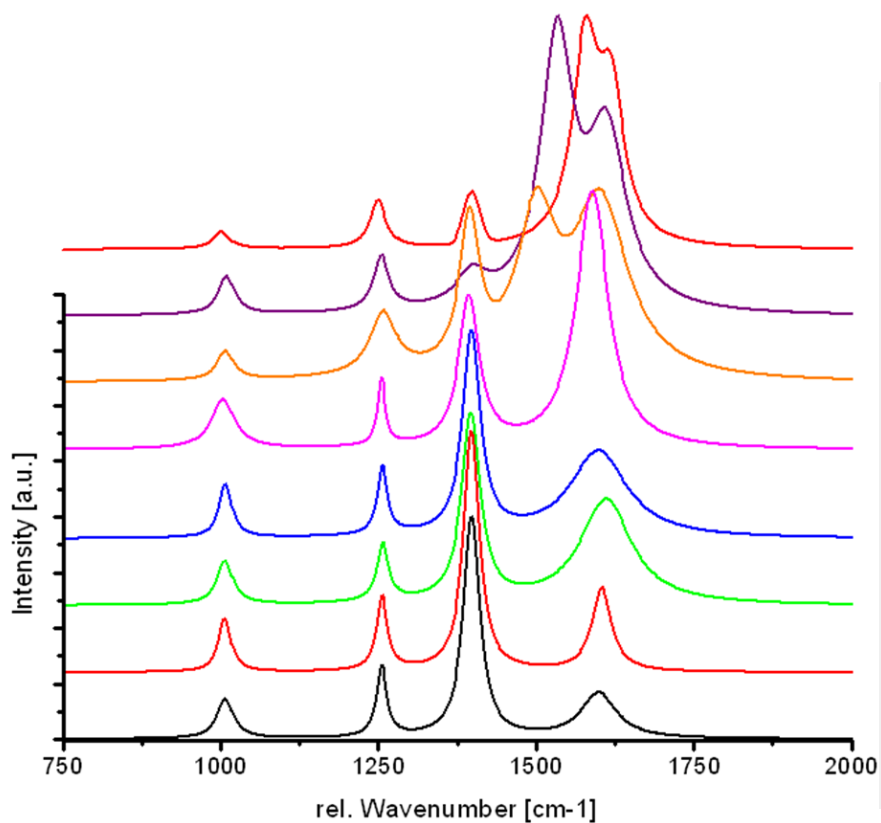


MBT(trans-Enol) on Au_{20}

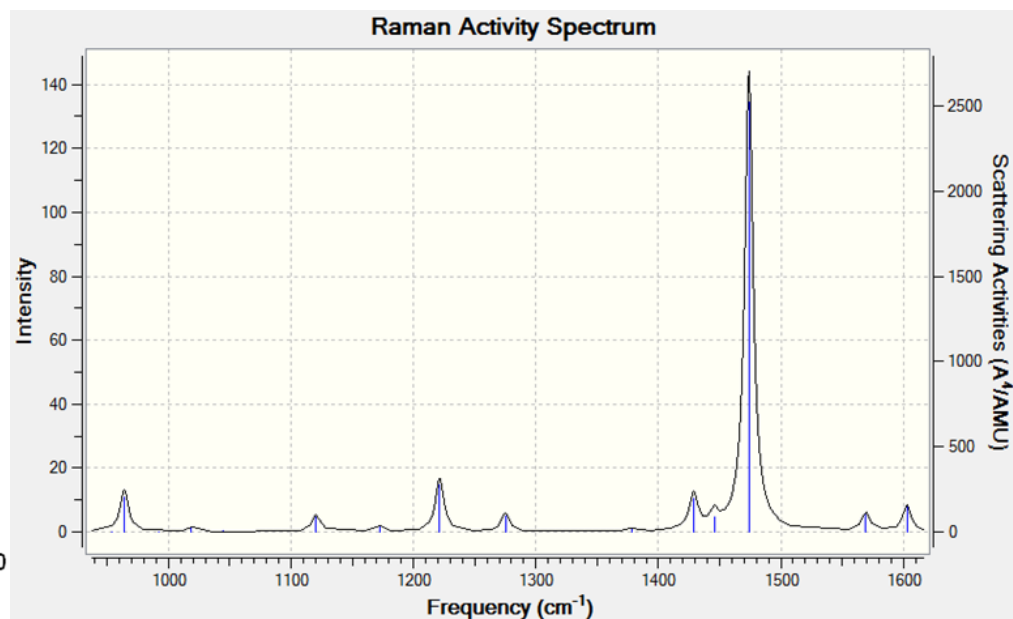
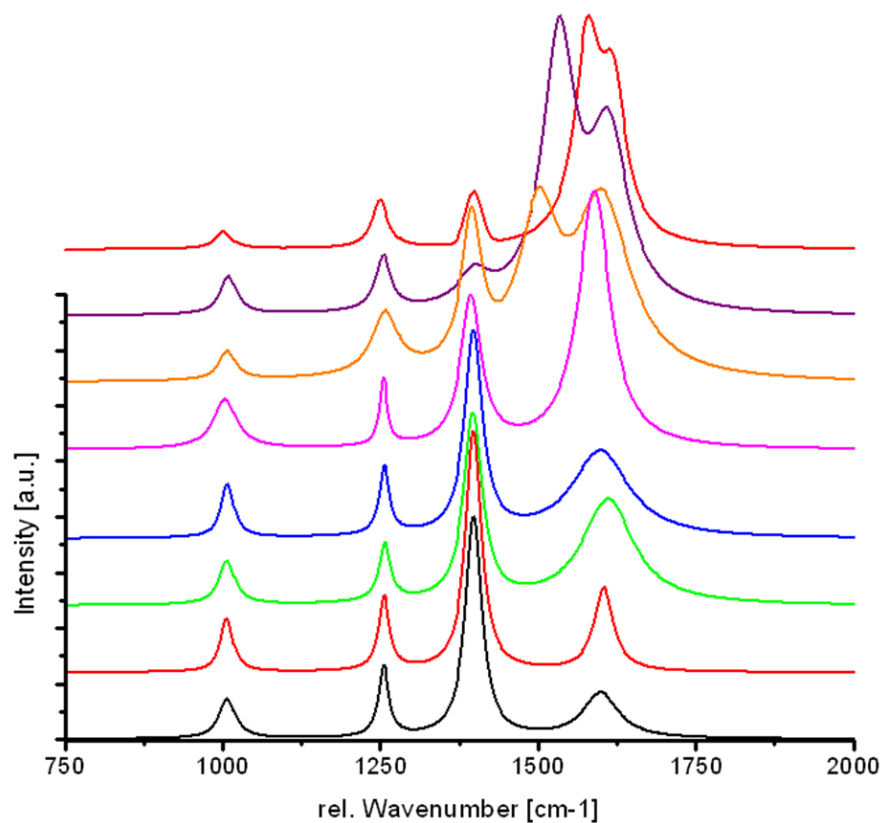


(D. Zhang, K. Braun)





MBT(Enolate) on Au₂₀



MBT(trans-Enol) on Au₂₀



- Lasse Jensen, Christine M. Aikens and George C. Schatz:
"Electronic structure methods for studying surface-enhanced Raman scattering"
Chem. Soc. Rev., 2008, **37**, 1061–1073.
- Hanning Chen, Jeffrey M. McMahon, Mark A. Ratner and George C. Schatz:
"Classical electrodynamics coupled to quantum mechanics for calculation of molecular optical properties: a RT-TDDFT/FDTD approach"
J. Chem. Phys. C, 2010, **114**, 14384–14392.
- Jonathan Mullin and George C. Schatz:
"Combined linear response quantum mechanics and classical electrodynamics (QM/ED) method for the calculation of surface-enhanced Raman spectra"
J. Phys. Chem. A, 2012, **116**, 1931–1938.